

# **University of Stuttgart**

Cluster of Excellence in Data-integrated Simulation Science

Project Coordinator: Prof. Dr. Andrea Barth Institute for Applied Analysis and Numerical Simulation

## Motivation: Sub-surface flow & Vehicular traffic



(Scalar) Conservation laws





#### uncertainties { We Roa

Weather data Road conditions

# Hyperbolic Conservation Laws with Random Discontinuous Fluxes Analysis and Simulation

Lukas Brencher

PN5

(2)

# Numerical approximation & Jump-adapted meshing

- Truncated Karhunen-Loève expansion for the Gaussian random field
- Depending on the specific construction of the jump field *P*: either exact evaluation possible or approximation via **Fourier inversion**.
- Finite Volume discretization
- Forward Euler scheme satifying the CFL stability condition
- Godunov flux:  $F_G(u, v) = \max\{f(\max(u, 0)), f(\min(v, 0))\}$

Jump-adapted wave-cell meshing



### Well-posedness

- Weak solutions not unique –> additional entropy condition necessary for uniqueness
- Discontinuous flux setting: infinitely many different entropy conditions
- Under suitable assumptions:
- Existence of pathwise entropy solution
- Uniqueness of pathwise entropy solution
- **X** Measurability of entropy solution fails with classical proofs

#### Theorem (Measurability of stochastic entropy solutions)

Let  $u_0 \in \mathcal{L}^q(\Omega; \mathcal{L}^p(\mathbb{R}^d))$ , with  $1 \leq q, p \leq \infty$ , be a stochastic initial condition to (1). Furthermore, for fixed  $\omega \in \Omega$ , assume that the solution  $u(\omega, \cdot, \cdot)$  takes values in a separable subspace  $S \subset \mathcal{L}^\infty$ . Then, the pathwise entropy solution to Problem (1) is strongly measurable in the sense that the mapping  $u : \Omega \to S$  is strongly measurable.

# Multiplicative flux & Stochastic jump coefficient



Cell interface at jump discontinuity (and domain boundary) Cell interfaces of additional wave-cells Cell interface resulting from (piecewise equidistant) refinement

# Pathwise solutions & convergence





#### **Multiplicative flux function**

 $f(\omega, t, \boldsymbol{x}, u) = \mathbf{a}(\omega, \boldsymbol{x})f(u)$ 

#### Stochastic jump coefficient

 $a(\omega, x) := \overline{a}(x) + \phi(W_{\mathcal{D}}(\omega, x)) + P(\omega, x)$ 

- $\overline{a} \in C(\mathbb{R}; \mathbb{R}_{\geq 0})$  is a deterministic, uniformly bounded mean function.
- $\phi \in C^1(\mathbb{R}; \mathbb{R}_{>0})$ . In our case:  $\phi(w) = \exp(w)$ .
- For a (zero-mean) Gaussian random field  $W \in L^2(\Omega; L^2(\mathbb{R}))$ associated to a non-negative, symmetric trace class (covariance) operator  $Q: L^2(\mathbb{R}) \to L^2(\mathbb{R})$ , the random field  $W_{\mathcal{D}} \in L^2(\Omega; L^2(\mathbb{R}))$  is defined as

 $W_{\mathcal{D}}(\omega, \mathbf{x}) = \begin{cases} W(\omega, \mathbf{x}), & \mathbf{x} \in \mathcal{D} \\ \min(W(\omega, \mathbf{x}), \sup_{\mathbf{x} \in \mathcal{D}} W(\omega, \mathbf{x})), & \mathbf{x} \in \mathbb{R} \setminus \mathcal{D} \end{cases}$ 

- *T* : Ω → B(D), ω ↦ {*T*<sub>1</sub>,...,*T*<sub>τ</sub>} is a random partition of D, i.e., the *T<sub>i</sub>* are disjoint open subsets of D with D
  = U<sup>τ</sup><sub>i=1</sub> T
  . The number of elements in T is a random variable τ : Ω → N on (Ω, A, P). For D<sub>I</sub> and D<sub>r</sub> being the left and right boundary of D, respectively, we define *T*<sub>0</sub> := (-∞, D<sub>I</sub>) and *T*<sub>τ+1</sub> := (D<sub>r</sub>, +∞).
- $(P_i, i \in \mathbb{N}_0)$  is a sequence of random variables on  $(\Omega, \mathcal{A}, \mathbb{P})$  with arbitrary non-negative distribution(s), which is independent of  $\tau$  (but not

#### References

- [1] Boris Andreianov, Kenneth H. Karlsen, and Nils H. Risebro. "A theory of L<sup>1</sup>-dissipative solvers for scalar conservation laws with discontinuous flux". In: Archive for Rational Mechanics and Analysis 201.1 (2011), pp. 27–86.
- [2] Emmanuel Audusse and Benoît Perthame. "Uniqueness for scalar conservation laws with discontinuous flux via adapted entropies". In: *Proceedings of the Royal Society of Edinburgh Section A: Mathematics* 135.2 (2005), pp. 253–265.

necessarily i.i.d.). Further we have

$$P: \Omega \times \mathcal{D} \to \mathbb{R}_{\geq 0}, \quad (\omega, x) \mapsto \sum_{i=0}^{\tau+1} \mathbf{1}_{\mathcal{T}_i}(x) P_i(\omega) .$$

### **Tools & Methods**

- Karhunen-Loève expansion
- Fourier inversion
- Finite Volume method
- Multilevel Monte Carlo method
- Forward Euler method
- Python / Matlab

- [3] Andrea Barth and Andreas Stein. "A study of elliptic partial differential equations with jump diffusion coefficients". In: SIAM/ASA Journal on Uncertainty Quantification 6.4 (2018), pp. 1707–1743.
- [4] Lukas Brencher and Andrea Barth. "Hyperbolic Conservation Laws with Stochastic Discontinuous Flux Functions". In: *International Conference on Finite Volumes for Complex Applications*. Springer, 2020, pp. 265–273.
- [5] Lukas Brencher and Andrea Barth. "Scalar conservation laws with stochastic discontinuous flux functions". In Preparation. 2021.
- [6] Lukas Brencher and Andrea Barth. "Stochastic conservation laws with discontinuous flux functions: The multidimensional case". In Preparation. 2021.
- [7] Gui-Qiang Chen, Nadine Even, and Christian Klingenberg. "Hyperbolic conservation laws with discontinuous fluxes and hydrodynamic limit for particle systems". In: *Journal of differential equations* 245.11 (2008), pp. 3095–3126.

# www.simtech.uni-stuttgart.de



